

## 1. MARCH 26

**1.1. Control Theory.** Like everything else, control theory is an ill-defined term, but we'll illustrate with an example.

Suppose  $\Omega$  is a smooth bounded domain. Let  $\Omega_T = \Omega \times (0, 1)$ , and consider the heat equation

$$(1.1) \quad \begin{aligned} (\partial_t - \Delta)u &= 0 \text{ in } \Omega_T \\ u(x, 0) &= u_0(x) \text{ on } \Omega. \end{aligned}$$

The question is, given a target function  $u_1(x)$  defined on  $\Omega$ , are there boundary values we can impose on  $\partial\Omega \times (0, 1)$ , such that the solution to (1.1) has the property that  $u(x, 1) = u_1(x)$ ?

In general the answer is no: because solutions to the heat equation are smooth, only certain  $u_1$  are acceptable. A better way to phrase the question might be to first define a space like

$$X = \{v_1 \in L^2(\Omega) \mid v_1(x) = v(x, 1), \text{ where } (\partial_t - \Delta)v = 0 \text{ in } \Omega \times (0, T), v_0(x) = v(x, 0) \in L^2(\Omega)\}.$$

Then we can ask, for a given  $v_1 \in X$ , is there  $f \in L^2(\partial\Omega \times (0, T))$  such that the solution to (1.1) with boundary condition  $u = g$  on  $\partial\Omega \times (0, T)$  satisfies  $u(x, T) = v_1(x)$ ?

Notice that if  $v_1 \in X$ , then by definition there exists  $v_0 \in L^2(\Omega)$  such that

$$\begin{aligned} (\partial_t - \Delta)v &= 0 \text{ in } \Omega \times (0, 1) \\ v(x, 0) &= v_0(x) \text{ on } \Omega \end{aligned}$$

and  $v(x, 1) = v_1(x)$ . Therefore if  $u$  solves (1.1) and  $u(x, 1) = v_1(x)$  then

$$\begin{aligned} (\partial_t - \Delta)(u - v) &= 0 \text{ in } \Omega \times (0, 1) \\ (u - v)(x, 0) &= (u_0 - v_0)(x) \text{ on } \Omega \end{aligned}$$

and  $u - v = 0$  at  $t = 1$ . The converse is true as well, so to answer the question of controllability, it suffices to understand the question of null-controllability.

**1.2. Null-controllability.** The basic question of null-controllability is the following: given a function  $u_0$  on  $\Omega$ , are there boundary conditions  $g \in L^2(\partial\Omega \times (0, 1))$  such that the equation

$$(1.2) \quad \begin{aligned} (\partial_t - \Delta)u &= 0 \text{ in } \Omega \times (0, 1) \\ u(x, 0) &= u_0(x) \text{ on } \Omega \\ u|_{\partial\Omega \times (0, 1)} &= g \end{aligned}$$

has the property  $u(x, 1) \equiv 0$ ? If so, we say that the initial condition  $u_0$  can be null controlled. More generally, we want to know what space  $E$  of functions defined on  $\Omega$  has the property that each  $u_0 \in E$  can be null controlled.

We say that a space  $E$  of functions on  $\Omega$  can be null-controlled if there exists  $C > 0$  such that for any  $u_0 \in E$ , there exists  $g \in L^2(\partial\Omega \times (0, 1))$  such that the solution  $u$  to the

boundary value problem

$$(1.3) \quad \begin{cases} (\partial_t - \Delta)u &= 0 \\ u(0, x) &= u_0(x) \\ u|_{\partial\Omega \times (0,1)} &= g \end{cases}$$

satisfies the condition  $u(1, x) \equiv 0$  and

$$(1.4) \quad \|g\|_{L^2(\partial\Omega \times (0,1))} \leq C \|u_0\|_E.$$

The last inequality is key to ensuring that the boundary conditions can be picked continuously, which is to say that the map from  $E$  to the necessary boundary condition is not too badly behaved.

To understand how on earth anyone solves a problem like this, we need to introduce the concept of the observability inequality.

### 1.3. Observability Inequality.

**Theorem 1.1.** *The following statements are equivalent:*

- *The space  $L^2(\Omega)$  is null-controlled.*
- *For all  $v \in H^2(\Omega_T)$  which solve  $(\partial_t + \Delta)v = 0$  on  $\Omega_T$  with  $v \equiv 0$  on  $\partial\Omega \times (0, 1)$ ,*

$$(1.5) \quad \|v\|_{L^2(\Omega \times \{0\})} \lesssim \|\partial_\nu v\|_{L^2(\partial\Omega \times (0,1))}.$$

*Proof.* First, suppose  $L^2(\Omega)$  is null-controlled, and suppose  $v \in H^2(\Omega_T)$  solves  $(\partial_t + \Delta)v = 0$  on  $\Omega_T$  with  $v \equiv 0$  on  $\partial\Omega \times (0, 1)$ . Now in particular  $v|_{t=0}$  is null-controlled, so we can set  $w$  to be a solution to  $(\partial_t - \Delta)w = 0$  with initial condition  $w(0, x) = v(0, x)$ , final condition  $w(1, x) \equiv 0$ , and boundary condition  $w|_{\partial\Omega \times (0,1)}$  satisfying (1.4). Then

$$\begin{aligned} 0 &= \int_{\Omega_T} w(\partial_t + \Delta)v \, dx \\ &= \int_{\Omega_T} (-\partial_t + \Delta)wv \, dx - \int_{\Omega} w(0, x)v(0, x) \, dx - \int_{\partial\Omega \times (0,1)} w\partial_\nu v \, dS \\ &= -\|v\|_{L^2(\Omega \times \{0\})}^2 - \int_{\partial\Omega \times (0,1)} w\partial_\nu v \, dS. \end{aligned}$$

Therefore

$$\|v\|_{L^2(\Omega \times \{0\})}^2 \leq C \|w\|_{L^2(\partial\Omega \times (0,1))} \|\partial_\nu v\|_{L^2(\partial\Omega \times (0,1))}.$$

By (1.4),

$$\|v\|_{L^2(\Omega \times \{0\})} \lesssim \|\partial_\nu v\|_{L^2(\partial\Omega \times (0,1))}.$$

Therefore (1.5) holds for each  $v \in H^2(\Omega_T)$  which solves  $(\partial_t + \Delta)v = 0$  on  $\Omega_T$  with  $v = 0$  on  $\partial\Omega \times (0, 1)$ .

Now suppose (1.5) holds for each  $v \in H^2(\Omega_T)$  which solves  $(\partial_t + \Delta)v = 0$  on  $\Omega_T$  with  $v = 0$  on  $\partial\Omega \times (0, 1)$ . Define

$$E = \{f \in L^2(\partial\Omega \times (0, 1)) \mid f = \partial_\nu v \text{ for some } v \in H^2(\Omega_T) \text{ s.t. } (\partial_t + \Delta)v = 0 \text{ and } v|_{\partial\Omega \times (0,1)} = 0.\}$$

Let  $u_0 \in L^2(\Omega)$ . Define the linear functional  $\varphi : E \rightarrow \mathbb{R}$  by

$$\varphi(f) = -(u_0, v|_{t=0})_\Omega.$$

Then (1.5) implies that

$$|\varphi(f)| \leq \|u_0\|_{L^2(\Omega)} \|v(0, x)\|_{L^2(\Omega)} \lesssim \|u_0\|_{L^2(\Omega)} \|f\|_{L^2(\partial\Omega \times (0,1))}.$$

Therefore  $\varphi$  is bounded as a map from  $E \subset L^2(\partial\Omega \times (0,1))$  to  $\mathbb{R}$ , and by Hahn-Banach, it has an extension to the whole space  $L^2(\partial\Omega \times (0,1))$ .

The extension is a linear functional on an  $L^2$  space, so by the Riesz representation theorem, there exists  $g \in L^2(\partial\Omega \times (0,1))$  with

$$\|g\|_{L^2(\partial\Omega \times (0,1))} = \|\varphi\| \lesssim \|u_0\|_{L^2(\Omega)}$$

and

$$\varphi(f) = (f, g)_{\partial\Omega \times (0,1)}$$

for all  $f \in L^2(\partial\Omega \times (0,1))$ . Therefore

$$-(u_0, v|_{t=0})_\Omega = (f, g)_{\partial\Omega \times (0,1)}.$$

Now let  $u$  be the unique solution to  $Pu = 0$  on  $\Omega_T$  with initial condition  $u(0, x) = u_0(x)$  and boundary conditions  $u|_{\partial\Omega \times (0,1)} = g$ . Then

$$-(u(0, x), v(0, x))_\Omega = (\partial_\nu v, u)_{\partial\Omega \times (0,1)}.$$

Meanwhile, an integration by parts gives

$$0 = ((-\partial_t + \Delta)u, v)_{\Omega_T} - (u, (\partial_t + \Delta)v)_{\Omega_T} = (u(1, x), v(1, x))_\Omega - (u(0, x), v(0, x))_\Omega - (\partial_\nu v, u)_{\partial\Omega \times (0,1)}.$$

Combining the above two equations gives

$$(u(1, x), v(1, x))_\Omega = 0.$$

This holds for any  $v$  solving the backwards heat equation  $(\partial_t + \Delta)v = 0$  on  $\Omega_T$  with  $v \equiv 0$  on  $\partial\Omega \times (0,1)$ , but for any given  $C^\infty(\Omega)$  function, we can solve the forwards heat equation on  $\Omega_T$  and then time reverse the solution to obtain any smooth final condition for  $v$  that we want. Therefore  $u(1, x) = 0$ , and it follows that  $L^2(\Omega)$  is null-controlled.  $\square$

So the problem of control comes down to an inequality, like any other self-respecting problem in analysis. On the other hand, this inequality does not look very much like any Carleman inequality we've seen so far.

In fact, trying to prove a Carleman inequality for the (adjoint) heat equation, we might end up with the following, very different looking result.

**Theorem 1.2.** *Suppose without loss of generality that  $\Omega$  does not contain 0 and is contained in the ball of radius  $R$  around 0. Let*

$$\varphi(x, t) = \frac{x^2 - R^2}{t(1-t)}.$$

Let

$$P_\varphi^* = e^{\tau\varphi}(\partial_t + \Delta)e^{-\tau\varphi}.$$

Then for all  $u \in H^2(\Omega_T)$  with  $u \equiv 0$  on  $\partial\Omega_T$ ,

$$\tau \|u\|_{L^2(\Omega_T)} \lesssim \|\sqrt{|\partial_\nu \varphi|} \partial_\nu u\|_{L^2(\partial\Omega \times (0,1))} + \|P_\varphi^* u\|_{L^2(\Omega_T)}.$$

This has an extra boundary term on the right side, compared to the Carleman estimates we're used to. But this isn't so strange – if we're not assuming that  $v$  vanishes to first order on the boundary, we should expect to pick up a boundary term from the integration by parts that we usually use to prove Carleman estimates. What's the relationship between the Carleman estimate and the observability inequality? I claim the Carleman estimate in fact proves the observability inequality, and hence controllability. Let's see why.

First, we have to make the standard substitution  $v = e^{-\tau\varphi}u$ , so we get

$$\tau \|e^{\tau\varphi}v\|_{L^2(\Omega_T)} \lesssim \|\sqrt{|\partial_\nu \varphi|} e^{\tau\varphi} \partial_\nu v\|_{L^2(\partial\Omega \times (0,1))} + \|e^{\tau\varphi}(\partial_t + \Delta)v\|_{L^2(\Omega_T)}$$

for all  $v$  such that  $e^{\tau\varphi}v \in H^2(\Omega_T)$  with  $v \equiv 0$  on  $\partial\Omega \times (0,1)$ . Now  $e^{\tau\varphi}$  and  $\sqrt{|\partial_\nu \varphi|} e^{\tau\varphi}$  is bounded above on  $\Omega_T$ , so in fact

$$\|e^{\tau\varphi}v\|_{L^2(\Omega_T)} \lesssim \|\partial_\nu v\|_{L^2(\partial\Omega \times (0,1))} + \|(\partial_t + \Delta)v\|_{L^2(\Omega_T)}$$

for all  $v \in H^2(\Omega_T)$  with  $v \equiv 0$  on  $\partial\Omega \times (0,1)$ . (Note that we don't have to specify that  $v$  is zero on the time boundaries  $t = 0$  or  $t = 1$ , because of the form of  $\varphi$  and the fact that  $v = e^{-\tau\varphi}u$ . Sadly  $e^{\tau\varphi}$  is not bounded below on  $\Omega_T$ . (why not?) But it is bounded below on  $\Omega \times (1/3, 2/3)$ , so

$$\|v\|_{L^2(\Omega \times (1/3, 2/3))} \lesssim \|\partial_\nu v\|_{L^2(\partial\Omega \times (0,1))} + \|(\partial_t + \Delta)v\|_{L^2(\Omega_T)}.$$

Now what? In the observability inequality, we make the crucial assumption that  $(\partial_t + \Delta)v = 0$ . If we apply this, we get

$$\|v\|_{L^2(\Omega \times (1/3, 2/3))} \lesssim \|\partial_\nu v\|_{L^2(\partial\Omega \times (0,1))}$$

for all  $v \in H^2(\Omega_T)$  with  $v \equiv 0$  on  $\partial\Omega \times (0,1)$  and  $(\partial_t + \Delta)v = 0$ . This is almost the observability inequality! We only have the wrong lower bound. But a simple energy argument fixes this: if  $(\partial_t + \Delta)v = 0$ , then

$$\begin{aligned} \partial_t \|v(x, t)\|_{L^2(\Omega)}^2 &= 2 \int_{\Omega} v(x, t) \partial_t v(x, t) dx \\ &= -2 \int_{\Omega} v(x, t) \Delta v(x, t) dx \\ &= 2 \|\nabla v(x, t)\|_{L^2(\Omega)}^2. \end{aligned}$$

Therefore  $\|v(x, t)\|_{L^2(\Omega)}^2$  is an increasing function of  $t$ , and

$$\|v\|_{L^2(\Omega \times \{0\})} \lesssim \|v\|_{L^2(\Omega \times (1/3, 2/3))}.$$

Therefore

$$\|v\|_{L^2(\Omega \times \{0\})} \lesssim \|\partial_\nu v\|_{L^2(\partial\Omega \times (0,1))}$$

for all  $v \in H^2(\Omega_T)$  with  $v \equiv 0$  on  $\partial\Omega \times (0,1)$  and  $(\partial_t + \Delta)v = 0$ . This proves the observability inequality, and hence the following theorem

**Theorem 1.3.** *Let  $\Omega$  be a smooth bounded domain. Then the space  $L^2(\Omega)$  is null-controlled for the heat equation.*

It only remains to prove the Carleman inequality.

*Proof of the Carleman inequality.* As usual, we begin by writing out in full

$$P_\varphi^* = \partial_t + \tau \partial_t \varphi + \Delta - \tau(\nabla \cdot \nabla \varphi + \nabla \varphi \cdot \nabla) + \tau^2 |\nabla \varphi|^2.$$

We write this in terms of symmetric and antisymmetric parts:

$$P_\varphi^* = A + B$$

where

$$A = \Delta + \tau^2 |\nabla \varphi|^2 + \tau \partial_t \varphi$$

and

$$B = \partial_t - \tau(\nabla \cdot \nabla \varphi + \nabla \varphi \cdot \nabla),$$

and we get

$$\|P_\varphi^* v\|_{L^2(\Omega_T)}^2 = \|Au\| + \|Bu\| + ([A, B]u, u) + \text{boundary terms}.$$

Because of the fact that  $u$  is zero on the boundary, only one boundary term in fact remains: it's the boundary term that you get from integrating by parts in the term

$$(2\tau \nabla \varphi \cdot \nabla u, \Delta u)_{\Omega_T} = 2\tau(\nabla \varphi \cdot \nabla u, \partial_\nu u)_{\partial\Omega \times (0,1)} + (\Delta 2\tau \nabla \varphi \cdot \nabla u, u)_{\Omega_T}$$

The tangential part of  $\nabla \varphi \cdot \nabla u$  is also irrelevant, because of the fact that  $u$  vanishes on the boundary, so in fact what we're left with is the boundary term

$$2\tau(\partial_\nu \varphi \partial_\nu u, \partial_\nu u)_{\partial\Omega \times (0,1)}.$$

This is bounded above by

$$C\tau \|\partial_\nu u\|_{L^2(\partial\Omega \times (0,1))}^2.$$

The commutator is bounded below as before (there are some extra terms, but disposing of them is merely an exercise in tedium). This finishes the proof of the Carleman estimate.  $\square$