

MA/PHY 507 Spring 2019  
Final Exam - 100 Points  
29 April 2019

INSTRUCTIONS: PLEASE WORK ALL FIVE PROBLEMS BELOW. NO BOOKS, PAPERS, OR NOTES ARE ALLOWED.

NAME: Solutions

PROBLEM	MAXIMUM GRADE	SCORE
1	25	
2	25	
3	25	
4	25	
TOTAL	100	

**Problem 1.** (25 points.) The heat equation for a temperature distribution  $u(x, t)$  takes the form

$$\frac{\partial u}{\partial t}(x, t) = D \frac{\partial^2 u}{\partial x^2}(x, t) - \lambda u(x, t),$$

for  $x \in \mathbb{R}$  and  $t \geq 0$ . The constants  $D$  and  $\lambda$  are both positive. Suppose the initial temperature distribution is  $u_0(x)$ . Use the Fourier transform to write the solution to the initial value problem as

$$u(x, t) = \int_{-\infty}^{\infty} K_t(x, y) u_0(y) dy.$$

Write a nice formula for the integral kernel  $K_t(x, y)$ .

**Useful identity:** The Fourier transform of  $e^{-ak^2}$ , for  $a > 0$ , is

$$\frac{1}{(2a)^{\frac{1}{2}}} e^{-\frac{x^2}{4a}}.$$

Fourier transform:  $\hat{u}(k, t) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-ikx} u(x, t) dx$

PDE becomes

$$\partial_t \hat{u}(k, t) = -Dk^2 \hat{u}(k, t) - \lambda \hat{u}(k, t) = -[Dk^2 + \lambda] \hat{u}(k, t)$$

Integrate:  $\hat{u}(k, t) = A e^{-(Dk^2 + \lambda)t}$

Initial condition:  $\hat{u}(k, t=0) = \hat{u}_0(k) = A$

$$\boxed{\hat{u}(k, t) = e^{-(Dk^2 + \lambda)t} \hat{u}_0(k)}$$

Take the inverse Fourier transform:

$$\begin{aligned} u(x, t) &= \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{ix \cdot k} \hat{u}(k, t) dk \\ &= \frac{1}{2\pi} \int_{\mathbb{R}} dk e^{i(x-y)k} e^{-(Dk^2 + \lambda)t} u_0(y) dy \end{aligned}$$

so  $K_t(x, y) = \frac{e^{-xt}}{2\pi} \int_{\mathbb{R}} dk e^{i(x-y)k} e^{-Dtk^2} dk$

$$= \left( \frac{e^{-xt}}{\sqrt{2\pi}} \right) \frac{1}{(2Dt)^{\frac{1}{2}}} e^{-\frac{|x-y|^2}{4Dt}}$$

so  $\boxed{K_t(x, y) = \frac{e^{-xt}}{(4\pi Dt)^{\frac{1}{2}}} e^{-\frac{|x-y|^2}{4Dt}}}$

Problem 2. (25 points.) Compute the integral:

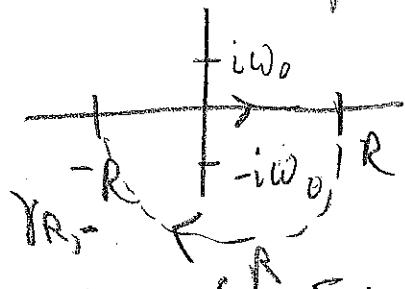
$$\int_{-\infty}^{\infty} \frac{e^{-i\omega t}}{\omega^2 + \omega_0^2} d\omega, \quad \omega_0 > 0,$$

for all real  $t \in \mathbb{R}$ . Carefully explain your procedure and choice of contour.

$F_t(\omega) = \frac{e^{-i\omega t}}{\omega^2 + \omega_0^2}$  is meromorphic with simple poles at  $\pm i\omega_0$

$$\text{Res}(F_t, i\omega_0) = \frac{e^{i\omega_0 t}}{2i\omega_0}$$

$$\text{Res}(F_t, -i\omega_0) = \frac{e^{-i\omega_0 t}}{-2i\omega_0}$$



$$\lim_{R \rightarrow \infty} \int_{-R}^R F_t(\omega) d\omega$$

(close the contours)  
 $t > 0$  LHP so  $e^{-i\omega t} = e^{-i(\text{Re}\omega)t} e^{(i\text{Im}\omega)t}$  decays exponentially for  $\text{Im}\omega < 0$ .

$$\oint F_t(\omega) d\omega = -2\pi i \text{Res}(F_t, -i\omega_0) = \frac{\pi}{\omega_0} e^{-\omega_0 t}$$

$t < 0$  UHP so  $e^{i\omega t} = e^{-i(\text{Re}\omega)t} e^{(i\text{Im}\omega)t}$  decays exponentially for  $\text{Im}\omega > 0$

$$\oint F_t(\omega) d\omega = 2\pi i \text{Res}(F_t, i\omega_0) = \frac{\pi}{\omega_0} e^{\omega_0 t}$$

$$\int_{-\infty}^{\infty} \frac{e^{-i\omega t}}{\omega^2 + \omega_0^2} d\omega = \frac{\pi}{\omega_0} e^{-|\omega_0|t}, \quad \text{Note: on the semi-circle } \text{Re}i\theta,$$

$\int_{\theta=0}^{\pi} |F_t(\omega)| \leq \frac{e^{-R\cos\theta|t|}}{((R e^{i\theta})^2 + \omega_0^2)} 2\pi R \rightarrow 0.$

**Problem 3.** (25 points.) Consider the heat equation  $\partial_t u = \Delta u$  in a sector of a disk described by  $0 \leq \rho \leq a$  and  $0 \leq \phi \leq \pi/2$  in two-dimensions. Suppose that the edges of the sector are kept at zero temperature.

- Write down the boundary conditions on  $u(\rho, \phi, t)$ .
- Find the most general real solution to the boundary-value problem obtained by separation of variables.

Useful information: Bessel's equation is:

$$x^2 y''(x) + xy'(x) + (\alpha^2 x^2 - \ell^2) y(x) = 0.$$

This ODE has  $J_\ell(\alpha x)$  as a regular solution. The Laplacian in polar coordinates is

$$\Delta = \rho^{-1} \partial_\rho \rho \partial_\rho + \rho^{-2} \partial_\phi^2.$$

Basic soln:  $\tilde{u}(\rho, \phi, t) = T(t) W(\rho, \phi)$ .

$$\frac{T'}{T} = -k^2 = \frac{\Delta W}{W} \Rightarrow T(t) = e^{-k^2 t} \text{ and } \Delta W = -k^2 W.$$

Radial PDE:  $W(\rho, \phi) = R(\rho) \Phi(\phi)$

$$\frac{1}{\rho R} \partial_\rho \rho \partial_\rho R + \frac{1}{\rho^2} \frac{\Phi''}{\Phi} = -k^2 \text{ or } \frac{k^2 R''}{R} + \rho R' + k^2 \rho^2 = -\frac{\Phi''}{\Phi}$$

Set  $\frac{\Phi''}{\Phi} = m^2$ ,  $m \in \mathbb{Z}$ , to insure single-valuedness;

$$\Phi(\phi) = A_m \cos m\phi + B_m \sin m\phi$$

Finally:  $\rho^2 R'' + \rho R' + [k^2 \rho^2 - m^2] R = 0$ .

Regular soln.  $R(\rho) = J_m(k\rho)$

$$\tilde{u}(t, \phi, \rho) = J_m(k\rho) (A_m \cos m\phi + B_m \sin m\phi) e^{-k^2 t}$$

BC:  $\tilde{u}(\rho=a, \phi, t) = 0$

$$\tilde{u}(\rho, \phi=0, t) = 0 = \tilde{u}(\rho, \phi=\frac{\pi}{2}, t)$$

$$\tilde{u}(\rho, 0, t) = A_m J_m(k\rho) e^{-k^2 t} = 0 \text{ so } A_m = 0$$

$$\tilde{u}(\rho, \frac{\pi}{2}, t) = B_m \sin(m\frac{\pi}{2}) J_m(k\rho) e^{-k^2 t} \text{ so } m = 2k, k \in \mathbb{N}.$$

$$\text{General Soln: } u(\rho, \phi, t) = \sum_{k=0}^{\infty} C_k \sin(2k\phi) J_{2k}(k\rho) e^{-k^2 t}$$

### Problem 3 (cont'd)

Finally BC:  $\hat{u}(r=a, \phi, t) = 0 \Leftrightarrow J_{2k}(\frac{k}{a}a) = 0$

$$\left(\frac{k}{a}\right) = \alpha_{2k,m}; \text{ mth zero of } J_{2k}$$

This introduces a second index  $m$ . The general soln is:

$$u(r, \phi, t) = \sum_{k=0}^{\infty} \sum_{m=1}^{\infty} C_{k,m} \sin(2k\phi) J_{2k}\left(\frac{\alpha_{2k,m}}{a} r\right) e^{-\frac{\alpha_{2k,m}}{a} t}$$

If one has an IC:  $u(r, \phi, t=0) = u_0(r, \phi)$

then:

$$u_0(r, \phi) = \sum_{k=0}^{\infty} \sum_{m=1}^{\infty} C_{k,m} \sin(2k\phi) J_{2k}\left(\frac{\alpha_{2k,m}}{a} r\right)$$

and the  $C_{k,m}$ 's are the Fourier-Bessel coefficients of  $u_0(r, \phi)$ .

**Problem 4.** (25 points.) The generating function  $g(x, t)$  for the Hermite polynomials  $H_n(x)$  is

$$g(x, t) = e^{-t^2+2xt} = \sum_{n=0}^{\infty} H_n(x) \frac{t^n}{n!}.$$

Prove the recursion relation for  $H_n(x)$  for  $n \geq 0$ :

$$2xH_n(x) - 2nH_{n-1}(x) = H_{n+1}(x).$$

Take  $\frac{\partial}{\partial t}$  of  $g(x, t)$  to avoid derivatives of  $H_n(x)$ :

$$\frac{\partial g(x, t)}{\partial t} = (2t + 2x)g(x, t) = \sum_{n=0}^{\infty} H_n(x) \frac{n!t^{n+1}}{n!}$$

Write the expansion of  $g$  and compare coeff. of  $t^k$  for all  $k$ :

$$\begin{aligned} (-2t + 2x)g(x, t) &= -2 \sum_{n=0}^{\infty} H_n(x) \frac{t^{n+1}}{n!} + 2 \sum_{n=0}^{\infty} xH_n(x) \frac{t^n}{n!} \\ &= \sum_{n=1}^{\infty} H_n(x) \frac{t^{n+1}}{(n-1)!} \end{aligned}$$

Reindex:

$$-2 \sum_{m=1}^{\infty} \frac{H_{m-1} t^m}{(m-1)!} + 2 \sum_{m=0}^{\infty} xH_m(x) \frac{t^m}{m!} = \sum_{m=0}^{\infty} H_{m+1}(x) \frac{t^m}{m!}$$

$$\sum_{m=0}^{\infty} \left[ -2mH_{m-1}(x) + 2xH_m(x) \right] \frac{t^m}{m!} = \sum_{m=0}^{\infty} H_{m+1}(x) \frac{t^m}{m!}$$

$$\text{So for } m \geq 0: -2mH_{m-1}(x) + 2xH_m(x) = H_{m+1}(x).$$