

Solutions to PS10

19.9.20 Clamped String on $[0, L]$
 $\partial_t^2 u = v^2 \partial_x^2 u$ with $u(x=0, t) = u(x=L, t) = 0$

Assume $u(x, t) = \sum b_n(t) \sin \frac{n\pi x}{L}$

Note: $\phi_n(x) = \sin \frac{n\pi x}{L}$ satisfies DBC

and $-\phi_n''(x) = \left(\frac{n\pi}{L}\right)^2 \phi_n(x)$

Substitute u into the WE

$$\partial_t^2 u = \sum_n b_n''(t) \phi_n(x) = -v^2 \sum_n \left(\frac{n\pi}{L}\right)^2 \phi_n(x) b_n(t)$$

$$b_n''(t) + v^2 \left(\frac{n\pi}{L}\right)^2 b_n(t) = 0.$$

$$b_n(t) = A_n \cos\left(\frac{v n \pi t}{L}\right) + B_n \sin\left(\frac{v n \pi t}{L}\right)$$

Now impose IVC:

$$u(x, 0) = f(x) \quad (\partial_t u)(x, 0) = g(x).$$

$$u(x, 0) = \sum_{n=1}^{\infty} A_n \sin \frac{n\pi x}{L} = f(x)$$

$$\Rightarrow A_n = \frac{2}{L} \int_0^L \sin\left(\frac{n\pi x}{L}\right) f(x) dx$$

$$\partial_t u(x, 0) = \sum_{n=1}^{\infty} \frac{n\pi v}{L} B_n \sin\left(\frac{n\pi x}{L}\right) \Rightarrow B_n = \frac{2}{n\pi v} \int_0^L \sin\left(\frac{n\pi x}{L}\right) g(x) dx$$

PS10-2

20.2.6 $\lim_{n \rightarrow \infty} \int f(x) \delta_n(x) dx = f(0)$.

Compute $\int f(x) \frac{1}{\sqrt{n}} e^{-n^2 x^2} dx$. Note: $\int \delta_n(x) dx = 1$

$u = nx$
 $= \frac{1}{\sqrt{n}} \int f\left(\frac{u}{n}\right) e^{-u^2} du$

For n large $f\left(\frac{u}{n}\right) = f(0) + \frac{u}{n} f'(0) + \frac{1}{2} \left(\frac{u}{n}\right)^2 f''(0) + \dots$

Leading term:

$$\int f(x) \delta_n(x) dx = f(0) \frac{1}{\sqrt{n}} \int e^{-u^2} du = f(0)$$

all other terms vanish as $n \rightarrow \infty$, like

$\bullet \frac{f'(0)}{\sqrt{n}} \frac{1}{n} \int e^{-u^2} u du = 0$ (symmetry)

$\bullet \frac{f''(0)}{2\sqrt{n}} \frac{1}{n^2} \int e^{-u^2} u^2 du \rightarrow 0$
 $n \rightarrow \infty$

Thus show that:

$$\lim_{n \rightarrow \infty} \int_{\mathbb{R}} f(x) \delta_n(x) dx = f(0) = \int_{\mathbb{R}} f(x) \delta(x) dx$$

As for the integral identity, note that

$$\begin{aligned}\hat{S}_n(k) &= \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} \delta_n(x) e^{-ikx} dx = \frac{n}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-n^2 x^2} e^{-ikx} dx \\ &= \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-u^2} e^{-i\frac{k}{n}u} du\end{aligned}$$

$$\text{and } \lim_{n \rightarrow \infty} \hat{S}_n(k) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-u^2} du = \frac{1}{\sqrt{2\pi}}$$

Taking the inverse FT

$$S(x) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{ikx} \frac{1}{\sqrt{2\pi}} dk = \frac{1}{2\pi} \int_{\mathbb{R}} e^{ikx} dk$$

$$= \frac{1}{2\pi} \int_{\mathbb{R}} e^{-ikt} dk$$

by change of $k \rightarrow -k$

20.3.6 Diffusion eqn:

$$-D \phi'' + K^2 D \phi = Q \delta(x)$$

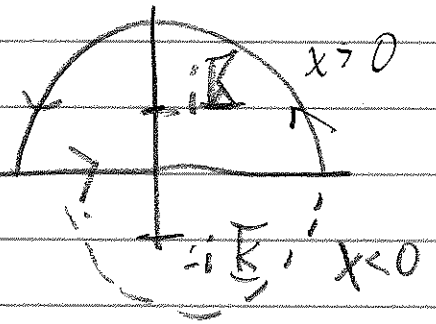
Take FT $\hat{\phi}(k) = \frac{1}{\sqrt{2\pi}} \int e^{-ikx} \phi(x) dx$

$$D k^2 \hat{\phi} + K^2 D \hat{\phi} = \frac{Q}{\sqrt{2\pi}}$$

$$\hat{\phi}(k) = \frac{1}{\sqrt{2\pi}} \frac{Q}{D(k^2 + K^2)} = \frac{1}{\sqrt{2\pi}} \frac{Q}{D} \frac{1}{k^2 + K^2}$$

$$\phi(x) = \frac{1}{\sqrt{2\pi}} \int e^{ikx} \left\{ \frac{1}{\sqrt{2\pi}} \left(\frac{1}{k^2 + K^2} \right) \right\} \left(\frac{Q}{D} \right)$$

Poles $k = \pm iK$, Take $\underline{K} > 0$



$x > 0$
 $\lim_{R \rightarrow \infty} \int_{-R}^R \frac{e^{ixk} dk}{(k - iK)(k + iK)} = 2\pi i \left(\frac{e^{-xK}}{2iK} \right) = \frac{\pi}{K} e^{-xK}$
 UHP

$x < 0$
 $\lim_{R \rightarrow \infty} \int_{-R}^R \frac{e^{ixk} dk}{(k - iK)(k + iK)} = -2\pi i \left(\frac{e^{xK}}{-2iK} \right) = \frac{\pi}{K} e^{xK}$
 LHP

$$\Rightarrow \boxed{\phi(x) = \frac{Q}{2KD} e^{-|x|K}}$$

ps 10.5

9.7.3

$$\partial_t \psi = a^2 \partial_x^2 \psi$$

$$\psi(x, t=0) = A \delta(x)$$

$$\begin{aligned}\psi(x, t) &= \int K_t(x, y) \psi(y, t=0) dy \\ &= A K_t(x, 0)\end{aligned}$$

Heat kernel at $y=0$.

$$\psi(x, t) = \frac{A}{(4a^2 t)^{1/2}} e^{-\frac{|x|^2}{4a^2 t}}$$